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## Site-Symmetry Restrictions on Thermal-Motion-Tensor Coefficients up to Rank 8

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### Abstract

An extension of the tables for the site-symmetry restrictions on the thermal-motion-tensor coefficients to include tensors up to rank 6 (up to rank 8 for cubic, hexagonal, tetragonal and trigonal site symmetries) is presented. The dependence relationships of these symmetric polar tensors are calculated in their natural crystallographic lattice frames using a direct method.

### Introduction

Recently there has been a growing interest in the study of anharmonicity in thermal motion; anharmonic parameters are essential, for example, in modelling the structures of fast ionic conductors or other disordered systems like ferroelectrics. These highly anharmonic systems very often need thermal parameters higher than fourth order to obtain a satisfactory description. Moreover, anisotropic modifications in a plane of hexagonal symmetry can be described only with sixth- (or higher) order terms. Nowadays the accuracy and precision of diffraction data are in general no longer an obstacle to a meaningful refinement of such high-order thermal parameters in a least-squares procedure, e.g. CsPbCl<sub>3</sub> (Hutton & Nelmes, 1981), Li<sub>3</sub>N (Zucker & Schulz, 1982), Ag<sub>3</sub>SI (Perenthaler & Schulz, 1981), PbF<sub>2</sub> (Schulz, Perenthaler & Zucker, 1982), RbAg<sub>4</sub>I<sub>5</sub> (Kuhns, 1983). Crystallographic program systems like *PROMETHEUS* (Zucker, Perenthaler, Kuhns, Bachmann & Schulz, 1983) allow for the routine insertion and refinement of parameters up to sixth order. To facilitate a general application, a tabulation of the symmetry constraints for all crystallographic site symmetries is clearly needed.

Whilst the symmetry restrictions of, for example, the elastic coefficients based on orthogonal lattice frames are at least partially available up to twelfth order (i.e. sixth-order elastic constants) (Chung & Li,

1974; Brendel, 1979; Fumi & Ripamonti, 1983), a tabulation for general (even- and odd-rank) symmetric polar tensors based on the natural crystallographic lattice frames is available only for tensors up to fourth order (Johnson & Levy, 1974). This paper presents the extension of that tabulation.

### Method

The method used to derive the dependence relations for symmetric polar tensor coefficients is based on the fundamental transformation law for tensors ('direct method'). A tensor coefficient remains unchanged when the transformation corresponds to the site symmetry:

$$C_T^{jklmnp\dots} = T_{jq}T_{kr}T_{ls}T_{mi}T_{nu}T_{pv}\dots C^{qrstuv\dots}$$

with  $C_T^{jklmnp\dots} = C^{qrstuv\dots}$  and  $1 \leq j, k, l, m, n, p, q, r, s, t, u, v \dots \leq 3$ . The point-group generators  $T$  were taken from Johnson & Levy (1974). For site symmetries having several generators all of them have been treated simultaneously. The homogeneous system of equations obtained (one equation for each coefficient) was reduced using a Gaussian elimination procedure. The order of assignment of independency is with increasing indices of the coefficients, except for the unmixed coefficients  $C^{jjjjj}$ , which have highest priority in every case.\* The correctness of the remaining system of equations giving the dependence relationships was checked algebraically by inserting numerical values. In addition, the number of unrestricted parameters was checked by comparison with results obtained by group-theoretical calculations (Sirotni, 1960).

\* This is the same ordering, as, for example, in the *PROMETHEUS* system (Zucker, Perenthaler, Kuhns, Bachmann & Schulz, 1983).

Table 1. Site-symmetry table giving key to tables 2A, B, for restrictions on the symmetry of various thermal-motion tensors (Hex denotes hexagonal axes)

Point symmetry at special position		Position $x, y, z$	Cross-reference number for		Point symmetry at special position		Position $x, y, z$	Cross-reference number for					
Symmetry axes	Point-group generators		2A	2B	Symmetry axes	Point-group generators		2A	2B				
$m\bar{3}m$		$4[0, 0, 1]$	$3[1, 1, 1]$	$\bar{1}$	0, 0, 0	E0	F1	$mm$	$2[1, 0, 0]$	$2[0, 0, 1]$	$x, y, z$	E44	F22
43m		$4[0, 0, 1]$	$3[1, 1, 1]$		0, 0, 0	E1	F1	$mm$	$2[1, 0, 0]$	$2[0, 1, 1]$	$x, 0, 0$	E45	F25
432		$4[0, 0, 1]$	$3[1, 1, 1]$		0, 0, 0	E0	F1	$mm$	$2[1, 1, 0]$	$2[0, 0, 1]$	$x, x, 0$	E46	F23
$m\bar{3}$		$3[1, 1, 1]$	$2[0, 0, 1]$	$\bar{1}$	0, 0, 0	E0	F2	$mm$	$2[1, \bar{1}, 0]$	$2[0, 0, 1]$	$x, \bar{x}, 0$	E47	F23
23		$3[1, 1, 1]$	$2[0, 0, 1]$		0, 0, 0	E1	F2	$mm$	$2[1, 0, 1]$	$2[0, 1, 0]$	$x, 0, x$	E48	F24
$6/mmm$	Hex	$6[0, 0, 1]$	$2[1, 0, 0]$	$\bar{1}$	0, 0, 0	E0	F3	$mm$	$2[1, 0, 1]$	$2[0, 1, 0]$	$x, 0, \bar{x}$	E49	F24
6m2	Hex	$6[0, 0, 1]$	$2[1, 0, 0]$		0, 0, 0	E5	F3	$mm$	$2[0, 1, 1]$	$2[1, 0, 0]$	$0, y, y$	E50	F25
6mm	Hex	$6[0, 0, 1]$	$2[1, 0, 0]$		0, 0, 0	E6	F3	$mm$	$2[0, 1, \bar{1}]$	$2[1, 0, 0]$	$0, y, \bar{y}$	E51	F25
622	Hex	$6[0, 0, 1]$	$2[1, 0, 0]$		0, 0, 0	E0	F4	$mm$	$2[0, 0, 1]$	$2[1, 0, 0]$	$0, 0, z$	E52	F26
$6/m$	Hex	$6[0, 0, 1]$	$\bar{1}$		0, 0, 0	E0	F4	$mm$	$2[0, 0, 1]$	$2[1, 1, 0]$	$0, 0, z$	E41	F23
6	Hex	$6[0, 0, 1]$			0, 0, 0	E24	F4	$mm$	$2[1, 0, 0]$	$2[0, 0, 1]$	$x, 0, 0$	E53	F27
6	Hex	$6[0, 0, 1]$			0, 0, z	E17	F4	$mm$	$2[2, 1, 0]$	$2[0, 0, 1]$	$2x, x, 0$	E54	F26
$4/mmm$		$4[0, 0, 1]$	$2[1, 0, 0]$	$\bar{1}$	0, 0, 0	E0	F5	$mm$	$2[1, 1, 0]$	$2[0, 0, 1]$	$x, x, 0$	E55	F27
$4/mmm$		$4[0, 1, 0]$	$2[0, 0, 1]$	$\bar{1}$	0, 0, 0	E0	F6	$mm$	$2[1, 2, 0]$	$2[0, 0, 1]$	$x, 2x, 0$	E56	F26
$4/mmm$		$4[1, 0, 0]$	$2[0, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F7	$mm$	$2[0, 1, 0]$	$2[0, 0, 1]$	$0, y, 0$	E57	F27
42m		$4[0, 0, 1]$	$2[1, 0, 0]$		0, 0, 0	E7	F5	$mm$	$2[1, \bar{1}, 0]$	$2[0, 0, 1]$	$x, \bar{x}, 0$	E47	F23
42m		$4[0, 0, 1]$	$2[1, 1, 0]$		0, 0, 0	E8	F5	222	$2[0, 0, 1]$	$2[1, 0, 0]$	$0, 0, 0$	E18	F22
42m		$4[0, 1, 0]$	$2[0, 0, 1]$		0, 0, 0	E9	F6	222	$2[0, 0, 1]$	$2[1, 1, 0]$	$0, 0, 0$	E19	F23
42m		$4[0, 1, 0]$	$2[1, 0, 1]$		0, 0, 0	E10	F6	222	$2[0, 1, 0]$	$2[1, 0, 1]$	$0, 0, 0$	E20	F24
42m		$4[1, 0, 0]$	$2[0, 1, 0]$		0, 0, 0	E11	F7	222	$2[1, 0, 0]$	$2[0, 1, 1]$	$0, 0, 0$	E21	F25
42m		$4[1, 0, 0]$	$2[0, 1, 1]$		0, 0, 0	E12	F7	222	$2[0, 0, 1]$	$2[1, 0, 0]$	$0, 0, 0$	E22	F26
4mm		$4[0, 0, 1]$	$2[1, 0, 0]$		0, 0, z	E25	F5	222	$2[0, 0, 1]$	$2[1, 1, 0]$	$0, 0, 0$	E19	F23
4mm		$4[0, 1, 0]$	$2[0, 0, 1]$		0, y, 0	E26	F6	222	$2[0, 0, 1]$	$2[0, 1, 0]$	$0, 0, 0$	E23	F27
4mm		$4[1, 0, 0]$	$2[0, 1, 0]$		$x, 0, 0$	E27	F7	2/m	$2[0, 0, 1]$	$\bar{1}$	$0, 0, 0$	E0	F28
422		$4[0, 0, 1]$	$2[1, 0, 0]$		0, 0, 0	E2	F5	2/m	$2[0, 1, 0]$	$\bar{1}$	$0, 0, 0$	E0	F29
422		$4[0, 1, 0]$	$2[0, 0, 1]$		0, 0, 0	E3	F6	2/m	$2[1, 0, 0]$	$\bar{1}$	$0, 0, 0$	E0	F30
422		$4[1, 0, 0]$	$2[0, 1, 0]$		0, 0, 0	E4	F7	2/m	$2[1, 1, 0]$	$\bar{1}$	$0, 0, 0$	E0	F31
4/m		$4[0, 0, 1]$	$\bar{1}$		0, 0, 0	E0	F14	2/m	$2[1, 1, 0]$	$\bar{1}$	$0, 0, 0$	E0	F32
4/m		$4[0, 1, 0]$	$\bar{1}$		0, 0, 0	E0	F15	2/m	$2[1, 0, 1]$	$\bar{1}$	$0, 0, 0$	E0	F33
4/m		$4[1, 0, 0]$	$\bar{1}$		0, 0, 0	E0	F16	2/m	$2[1, 0, \bar{1}]$	$\bar{1}$	$0, 0, 0$	E0	F34
4		$4[0, 0, 1]$			0, 0, 0	E28	F14	2/m	$2[0, 1, 1]$	$\bar{1}$	$0, 0, 0$	E0	F35
4		$4[0, 1, 0]$			0, 0, 0	E29	F15	2/m	$2[0, 1, \bar{1}]$	$\bar{1}$	$0, 0, 0$	E0	F36
4		$4[1, 0, 0]$			0, 0, 0	E30	F16	2/m	$2[0, 0, 1]$	$\bar{1}$	$0, 0, 0$	E0	F28
4		$4[0, 0, 1]$			0, 0, z	E31	F14	2/m	$2[1, 0, 0]$	$\bar{1}$	$0, 0, 0$	E0	F37
4		$4[0, 1, 0]$			0, y, 0	E32	F15	2/m	$2[2, 1, 0]$	$\bar{1}$	$0, 0, 0$	E0	F38
4		$4[1, 0, 0]$			$x, 0, 0$	E33	F16	2/m	$2[1, 1, 0]$	$\bar{1}$	$0, 0, 0$	E0	F31
3m		$3[1, 1, 1]$	$2[1, \bar{1}, 0]$	$\bar{1}$	0, 0, 0	E0	F8	2/m	$2[1, 2, 0]$	$\bar{1}$	$0, 0, 0$	E0	F39
3m		$3[1, 1, \bar{1}]$	$2[1, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F9	2/m	$2[0, 1, 0]$	$\bar{1}$	$0, 0, 0$	E0	F40
3m		$3[1, \bar{1}, 1]$	$2[1, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F10	2/m	$2[1, \bar{1}, 0]$	$\bar{1}$	$0, 0, 0$	E0	F32
3m		$3[\bar{1}, 1, 1]$	$2[1, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F11	m	$2[0, 0, 1]$		$x, y, 0$	E76	F28
3m	Hex	$3[0, 0, 1]$	$2[1, 0, 0]$	$\bar{1}$	0, 0, 0	E0	F12	m	$2[0, 1, 0]$		$x, 0, z$	E77	F29
3m	Hex	$3[0, 0, 1]$	$2[1, 2, 0]$	$\bar{1}$	0, 0, 0	E0	F13	m	$2[1, 0, 0]$		$0, y, z$	E78	F30
3m		$3[1, 1, 1]$	$2[1, \bar{1}, 0]$		$x, x, x$	E34	F8	m	$2[1, 1, 0]$		$x, \bar{x}, z$	E79	F31
3m		$3[1, 1, \bar{1}]$	$2[1, \bar{1}, 0]$		$x, x, \bar{x}$	E35	F9	m	$2[1, \bar{1}, 0]$		$x, x, z$	E80	F32
3m		$3[1, \bar{1}, 1]$	$2[1, 1, 0]$		$x, \bar{x}, x$	E36	F10	m	$2[1, 0, 1]$		$x, y, \bar{x}$	E81	F33
3m		$3[\bar{1}, 1, 1]$	$2[1, 1, 0]$		$\bar{x}, x, x$	E37	F11	m	$2[1, 0, \bar{1}]$		$x, y, x$	E82	F34
3m	Hex	$3[0, 0, 1]$	$2[1, 0, 0]$		0, 0, z	E38	F12	m	$2[0, 1, 1]$		$x, y, y$	E83	F35
3m	Hex	$3[0, 0, 1]$	$2[1, 2, 0]$		0, 0, z	E39	F13	m	$2[0, 1, \bar{1}]$		$x, y, \bar{y}$	E84	F36
32		$3[1, 1, 1]$	$2[1, 1, 0]$		0, 0, 0	E13	F8	m	$2[0, 0, 1]$		$x, y, 0$	E76	F28
32		$3[1, 1, \bar{1}]$	$2[1, 1, 0]$		0, 0, 0	E14	F9	m	$2[1, 0, 0]$		$x, 2x, z$	E85	F37
32		$3[1, \bar{1}, 1]$	$2[1, 1, 0]$		0, 0, 0	E15	F10	m	$2[2, 1, 0]$		$0, y, z$	E86	F38
32		$3[\bar{1}, 1, 1]$	$2[1, 1, 0]$		0, 0, 0	E16	F11	m	$2[1, 1, 0]$		$x, \bar{x}, z$	E79	F31
32	Hex	$3[0, 0, 1]$	$2[1, 0, 0]$		0, 0, 0	E5	F12	m	$2[1, 2, 0]$		$x, 0, z$	E87	F39
32	Hex	$3[0, 0, 1]$	$2[1, 2, 0]$		0, 0, 0	E6	F13	m	$2[0, 1, 0]$		$2x, x, z$	E88	F40
3		$3[1, 1, 1]$			0, 0, 0	E0	F17	m	$2[1, \bar{1}, 0]$		$x, x, z$	E80	F32
3		$3[1, 1, \bar{1}]$			0, 0, 0	E0	F18	2	$2[0, 0, 1]$		$0, 0, z$	E63	F28
3		$3[1, \bar{1}, 1]$			0, 0, 0	E0	F19	2	$2[0, 1, 0]$		$0, y, 0$	E64	F29
3		$3[\bar{1}, 1, 1]$			0, 0, 0	E0	F20	2	$2[1, 0, 0]$		$x, 0, 0$	E65	F30
3	Hex	$3[0, 0, 1]$			0, 0, 0	E0	F21	2	$2[1, 1, 0]$		$x, x, 0$	E66	F31
3		$3[1, 1, 1]$			$x, x, x$	E58	F17	2	$2[1, \bar{1}, 0]$		$x, \bar{x}, 0$	E67	F32
3		$3[1, 1, \bar{1}]$			$x, x, \bar{x}$	E59	F18	2	$2[1, 0, 1]$		$x, 0, x$	E68	F33
3		$3[1, \bar{1}, 1]$			$x, \bar{x}, x$	E60	F19	2	$2[1, 0, \bar{1}]$		$x, 0, \bar{x}$	E69	F34
3		$3[\bar{1}, 1, 1]$			$\bar{x}, x, x$	E61	F20	2	$2[0, 1, 1]$		$0, y, y$	E70	F35
3	Hex	$3[0, 0, 1]$			0, 0, z	E62	F21	2	$2[0, 1, \bar{1}]$		$0, y, \bar{y}$	E71	F36
mmm		$2[0, 0, 1]$	$2[1, 0, 0]$	$\bar{1}$	0, 0, 0	E0	F22	2	$2[0, 0, 1]$		$0, 0, z$	E63	F28
mmm		$2[0, 0, 1]$	$2[1, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F23	2	$2[1, 0, 0]$		$x, 0, 0$	E72	F37
mmm		$2[0, 1, 0]$	$2[1, 0, 1]$	$\bar{1}$	0, 0, 0	E0	F24	2	$2[2, 1, 0]$		$2x, x, 0$	E73	F38
mmm		$2[1, 0, 0]$	$2[0, 1, 1]$	$\bar{1}$	0, 0, 0	E0	F25	2	$2[1, 1, 0]$		$x, x, 0$	E66	F31
mmm	Hex	$2[0, 0, 1]$	$2[1, 0, 0]$	$\bar{1}$	0, 0, 0	E0	F26	2	$2[1, 2, 0]$		$x, 2x, 0$	E74	F39
mmm	Hex	$2[0, 0, 1]$	$2[1, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F23	2	$2[0, 1, 0]$		$0, y, 0$	E75	F40
mmm	Hex	$2[0, 0, 1]$	$2[0, 1, 0]$	$\bar{1}$	0, 0, 0	E0	F27	2	$2[1, 1, 0]$		$x, \bar{x}, 0$	E67	F32
mm		$2[0, 0, 1]$	$2[1, 0, 0]$		0, 0, z	E40	F22	$\bar{1}$	$\bar{1}$		$0, 0, 0$	E0	F41
mm		$2[0, 0, 1]$	$2[1, 1, 0]$		0, 0, z	E41	F23	$\bar{1}$	$\bar{1}$		$0, 0, 0$	E0	F41
mm		$2[0, 1, 0]$	$2[0, 0, 1]$		0, y, 0	E42	F22	1	1		$x, y, z$	E89	F41
mm		$2[0, 1, 0]$	$2[1, 0, 1]$		0, y, 0	E43	F24	1	Hex	1	$x, y, z$	E89	F41

Table 2A. Symmetry restrictions on coefficients in fifth-rank symmetric polar tensors

	A	B	C	D	E	F	G	H	I	J	K	L	M	N	P	Q	R	S	T	U	V		
	1	2	3	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	2	2	2	2	
Number of independent coefficients	1	2	3	1	1	1	1	1	1	1	1	1	2	2	2	2	3	2	2	2	3	3	
Cross reference	1	2	3	1	1	2	2	3	2	2	3	3	2	2	3	3	3	2	3	3	3	3	
E0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	
E1	1	0	0	0	0	0	0	G	0	0	0	0	0	G	0	G	0	0	0	0	0	0	
E2	1	0	0	0	0	0	0	G	0	0	0	0	0	-G	0	0	0	0	0	0	0	0	
E3	1	0	0	0	0	0	0	G	0	0	0	0	0	0	0	-G	0	0	0	0	0	0	
E4	1	0	0	0	0	0	0	0	0	0	0	0	0	N	0	-N	0	0	0	0	0	0	
E5	2	0	0	0	D	0	D	0	D	0	K	0	D	0	K	0	0	0	0	0	0	0	
E6	2	A	-A	0	A/2	0	A/10	0	H	-A/10	0	H/2	0	-A/2	0	-H/2	0	0	0	0	0	-H	
E7	2	0	0	0	0	0	0	G	0	0	0	0	0	G	0	Q	0	0	0	0	0	0	
E8	2	0	0	0	0	E	0	0	0	0	0	L	0	0	0	0	0	0	-E	0	-L	0	
E9	2	0	0	0	0	0	0	G	0	0	0	0	0	N	0	G	0	0	0	0	0	0	
E10	2	0	0	0	D	0	0	0	I	0	0	0	0	0	0	0	0	0	0	-I	0	-D	
E11	2	0	0	0	0	0	0	G	0	0	0	0	0	N	0	N	0	0	0	0	0	0	
E12	2	0	0	0	0	0	0	F	0	-F	0	0	0	M	0	0	0	-M	0	0	0	0	
E13	2	0	0	0	D	-D	F	0	-F	-F	0	0	F	-D	0	0	0	D	D	F	F	-D	
E14	2	0	0	0	D	D	F	0	-F	-F	0	0	-F	-D	0	0	0	D	-D	F	F	-D	
E15	2	0	0	0	D	D	F	0	-F	F	0	0	F	D	0	0	0	-D	-D	-F	-F	-D	
E16	2	0	0	0	D	-D	F	0	-F	F	0	0	-F	D	0	0	0	-D	D	-F	F	-D	
E17	3	0	0	0	C	0	E	0	E/2	0	0	E/2	0	L	0	E/2	0	L/2	0	E	0	L	0
E18	3	0	0	0	0	0	0	G	0	0	0	0	0	N	0	Q	0	0	0	0	0	0	
E19	3	0	0	0	0	E	0	G	0	0	0	L	0	-G	0	0	0	-E	0	-L	0	0	
E20	3	0	0	0	D	0	0	G	0	I	0	0	0	0	0	-G	0	0	0	-I	0	-D	
E21	3	0	0	0	0	0	F	0	-F	0	0	0	0	M	N	0	-M	0	0	0	0	0	
E22	3	0	0	0	0	0	0	G	0	0	2G	0	L	0	2G	0	L	0	S	0	0	0	
E23	3	0	0	0	0	0	0	G	0	0	G	0	0	N	0	Q	0	14*	0	Q	0	0	
E24	4	A	-A	0	D	0	1*	0	H	3*	0	K	0	7*	0	13*	0	0	0	-H	0	0	
E25	4	0	0	0	C	0	E	0	0	0	J	0	L	0	0	0	0	0	E	0	L	0	
E26	4	0	0	0	B	0	D	0	0	0	I	0	K	0	0	0	0	0	0	I	0	D	
E27	4	A	0	0	0	0	F	0	F	0	0	0	0	M	0	P	0	M	0	0	0	0	
E28	4	0	0	0	0	E	0	G	0	0	0	L	0	G	0	Q	0	-E	0	-L	0	0	
E29	4	0	0	0	D	0	0	G	0	I	0	0	0	N	0	G	0	0	-I	0	-D	0	
E30	4	0	0	0	0	F	G	-F	0	0	0	0	0	M	N	0	-M	0	0	0	0	0	
E31	5	0	0	0	C	0	E	0	G	0	J	0	L	0	-G	0	0	0	E	0	L	0	
E32	5	0	0	0	B	0	D	0	0	G	0	I	0	K	0	0	0	-G	0	I	0	D	
E33	5	A	0	0	0	0	0	F	0	F	0	0	0	M	N	P	-N	M	0	0	0	0	
E34	5	A	A	A	D	D	F	G	F	F	J	J	F	D	G	J	G	D	D	F	F	D	
E35	5	A	A	-A	D	D	-D	F	G	F	F	J	-J	-F	D	G	-J	G	D	-D	F	D	
E36	5	A	-A	A	D	-D	F	G	F	-F	J	-J	F	-D	G	J	G	-D	-D	-F	F	D	
E37	5	A	-A	-A	D	D	F	G	F	-F	J	J	-F	-D	G	-J	G	-D	D	-F	-F	D	
E38	5	A	-A	C	A/2	E	A/10	E/2	H	-A/10	E/2	H/2	L	-A/2	E/2	-H/2	L/2	0	E	-H	L	0	
E39	5	0	0	0	C	D	E	D	E/2	0	D	E/2	K	L	D	E/2	K	L/2	0	E	0	L	0
E40	6	0	0	0	C	0	E	0	0	0	0	J	0	L	0	0	0	0	0	S	0	U	0
E41	6	0	0	0	C	0	E	0	G	0	0	J	0	L	0	G	0	Q	0	E	0	L	0
E42	6	0	0	0	B	0	D	0	0	0	I	0	K	0	0	0	0	0	0	0	T	0	V
E43	6	0	0	0	B	0	D	0	0	G	0	I	0	K	0	0	N	0	C	0	I	0	D
E44	6	A	0	0	0	0	F	0	H	0	0	0	0	M	0	P	0	R	0	0	0	0	0
E45	6	A	0	0	0	0	F	G	H	0	0	0	0	M	N	P	N	M	0	0	0	0	0
E46	6	A	-A	0	D	0	F	0	H	F	0	K	0	D	0	K	0	R	0	H	0	R	
E47	6	A	-A	0	D	0	F	0	H	-F	0	K	0	-D	0	-K	0	R	0	-H	0	-R	
E48	6	A	0	0	E	F	0	H	0	J	0	H	M	0	J	0	E	M	0	F	0	0	
E49	6	A	0	-A	0	E	F	0	H	0	G	0	-H	M	0	-J	0	-E	-M	0	-F	0	
E50	6	0	0	B	B	D	D	0	0	0	I	J	I	0	0	0	0	0	S	T	T	S	
E51	6	0	0	B	-B	D	-D	0	0	0	I	J	-I	0	0	0	0	0	S	T	-T	-S	
E52	6	0	0	-C	0	E	0	G	0	0	J	0	L	0	10*	0	Q	0	15*	0	2Q	0	
E53	6	0	0	0	C	0	E	0	E/2	0	0	J	0	L	0	11*	0	L/2	0	S	0	U	0
E54	6	A	0	0	D	0	F	0	H	4*	0	K	0	4*	0	K	0	R	0	0	0	0	
E55	6	A	B	0	A/2	0	F	0	H	5*	0	H/2	0	8*	0	P	0	R	0	16*	0	R/2	
E56	6	A	B	0	D	0	2*	0	H	6*	0	K	0	B/2	0	12*	0	R	0	17*	0	2R	
E57	6	0	0	B	0	D	0	0	0	I	0	K	0	9*	0	K	0	0	0	T	0	V	
E58	7	A	A	A	D	E	F	G	H	H	J	J	F	E	G	J	G	D	D	F	H	E	
E59	7	A	A	-A	D	E	F	G	H	H	J	-J	-F	-E	G	-J	G	D	-D	F	-H	-E	
E60	7	A	-A	A	D	E	F	G	H	-H	J	-J	F	E	G	J	G	-D	-D	-F	H	-E	
E61	7	A	-A	-A	D	E	F	G	H	-H	J	J	-F	-E	G	-J	G	-D	D	-F	-H	E	
E62	7	A	-A	C	D	E	1*	E/2	H	3*	E/2	K	L	7*	E/2	13*	L/2	0	E	-H	L	0	
E63	9	0	0	0	C	0	E	0	G	0	0	J	0	L	0	N	0	Q	0	S	0	U	0
E64	9	0	0	0	B	0	D	0	0	G	0	I	0	K	0	0	N	0	Q	0	T	0	V
E65	9	A	0	0	0	0	F	G	H	0	0	0	0	M	N	P	Q	R	0	0	0	0	
E66	9	A	A	0	D	E	F	G	H	F	0	K	L	D	-G	K	0	R	0	-E	H	-L	R
E67	9	A	-A	0	D	E	F	G	H	-F	0	K	L	-D	-G	-K	0	R	0	-E	H	-L	R
E68	9	A	0	0	A	D	E	F	G	H	I	J	0	H	M	0	J	-G	E	M	-I	F	-D
E69	9	A	0	-A	D	E	F	G	H	I	J	0	-H	M	0	-J	-G	-E	-M	-I	-F	-D	
E70	9	0	0	B	B	D	D	F															

Table 2A (cont.)

		A	B	C	D	E	F	G	H	I	J	K	L	M	N	P	Q	R	S	T	U	V	
	Number of independent coefficients	1	2	3	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	2	2	2	2
	Cross reference	1	2	3	1	1	1	1	1	2	2	2	3	2	2	2	3	3	3	2	3	3	3
E81	12	A	B	-A	D	E	F	G	H	I	J	K	-H	M	N	-J	G	-E	-M	I	-F	D	
E82	12	A	B	A	D	E	F	G	H	I	J	K	H	M	N	J	G	E	M	I	F	D	
E83	12	A	B	-B	D	-D	F	G	F	I	J	-J	-I	M	N	P	N	M	S	T	-T	-S	
E84	12	A	B	B	D	D	F	G	F	I	J	J	I	M	N	P	N	M	S	T	T	S	
E85	12	A	B	C	D	E	2*	G	H	6*	J	K	L	B/2	10*	12*	N	M	15*	17*	2N	2M	
E86	12	0	B	C	D	E	D	E/2	0	I	J	K	L	2I	11*	K	L/2	0	S	T	U	V	
E87	12	A	0	C	D	E	F	G	H	4*	J	K	L	4*	10*	K	N	M	15*	0	2N	0	
E88	12	A	B	C	A/2	E	F	E/2	H	5*	J	H/2	L	8*	11*	P	L/2	R	S	16*	U	U/2	
E89	21	A	B	C	D	E	F	G	H	I	J	K	L	M	N	P	Q	R	S	T	U	V	

\* Notes, 1 -2A/5+D; 2 -3A/5+B/10+3D/2; 3 -3A/5+D; 4 -D+2F; 5 -A/4+3F/2; 6 -2A/5+B/5+D; 7 -A+D; 8 -A/5+2B/5+F; 9 -D+2J; 10 -2G+3J; 11 -E/4+3J/2; 12 -2H+3K; 13 -H+K; 14 -G+2N; 15 -4G+6J; 16 -H/4+3P/2; 17 -4H+6K.

Table 2B. Symmetry restrictions on coefficients in sixth-rank symmetric polar tensors

		A	B	C	D	E	F	G	H	I	J	K	L	M	N	P	Q	R	S	T	U	V	W	X	Y	Z	a	b	c	
	Number of independent parameters	1	2	3	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	2	2	2	2	
	Cross reference	1	2	3	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	2	2	2	2	
F1	3	A	A	A	0	0	F	0	F	0	0	0	0	F	0	P	0	F	0	0	0	0	0	0	0	F	0	F	0	
F2	4	A	A	A	0	0	F	0	H	0	0	0	0	H	0	P	0	F	0	0	0	0	0	0	0	F	0	H	0	
F3	5	A	A	C	A/2	0	F	0	H	1*	0	H/2	0	F	0	H/2	0	R	A/2	0	H/2	0	R/2	0	0	H	0	R	0	
F4	6	A	A	C	D	0	F	0	H	2*	0	H/2	0	5*	0	H/2	0	R	11*	0	H/2	0	R/2	0	0	H	0	R	0	
F5	6	A	A	C	0	0	F	0	H	0	0	0	0	F	0	P	0	R	0	0	0	0	0	0	0	H	0	R	0	
F6	6	A	B	A	0	0	F	0	H	0	0	0	0	M	0	P	0	H	0	0	0	0	0	0	0	M	0	F	0	
F7	6	A	B	B	0	0	F	0	F	0	0	0	0	M	0	P	0	M	0	0	0	0	0	0	0	0	Z	0	Z	0
F8	7	A	A	A	D	D	F	G	F	I	J	J	I	F	J	P	J	F	D	G	J	J	G	D	D	F	I	F	D	
F9	7	A	A	A	D	-D	F	G	F	I	J	-J	-I	F	J	P	J	F	D	G	-J	J	-G	-D	-D	F	-I	F	-D	
F10	7	A	A	A	D	-D	F	G	F	I	J	-J	-I	F	-J	P	-J	F	D	-G	-J	-J	G	-D	D	F	I	F	D	
F11	7	A	A	A	D	D	F	G	F	I	J	J	I	F	-J	P	-J	F	D	-G	J	-J	-G	D	-D	F	-I	F	-D	
F12	7	A	A	C	A/2	E	F	E/2	H	1*	E/10	H/2	I	F	-E/10	H/2	1/2	R	A/2	-E/2	H/2	-1/2	R/2	0	-E	H	-I	R	0	
F13	7	A	A	C	A/2	0	F	G	H	1*	G	H/2	0	F	G	H/2	Q	R	A/2	G	H/2	Q	R/2	0	0	H	0	R	0	
F14	8	A	A	C	D	0	F	0	H	0	0	K	0	F	0	P	0	R	-D	0	0	0	0	0	0	H	0	R	0	
F15	8	A	B	A	0	E	F	0	H	0	J	0	0	M	0	P	0	H	0	0	0	-J	0	-E	0	M	0	F	0	
F16	8	A	B	B	0	0	F	0	H	0	0	0	0	M	N	P	-N	M	0	0	0	0	0	0	0	Y	Z	0	Z	-Y
F17	10	A	A	A	D	E	F	G	H	I	J	K	I	H	K	P	J	F	E	G	J	K	G	D	D	F	I	H	E	
F18	10	A	A	A	D	E	F	G	H	I	J	K	-I	H	-K	P	J	F	-E	G	-J	-K	-G	-D	-D	F	-I	H	E	
F19	10	A	A	A	D	E	F	G	H	I	J	K	-I	H	K	P	-J	F	-E	-G	-J	-K	G	-D	D	F	I	H	-E	
F20	10	A	A	A	D	E	F	G	H	I	J	K	I	H	-K	P	-J	F	E	-G	J	K	-G	D	-D	F	-I	H	-E	
F21	10	A	A	C	D	E	F	G	H	2*	4*	H/2	L	5*	7*	H/2	Q	R	11*	13*	H/2	18*	R/2	0	-E	H	-L	R	0	
F22	10	A	B	C	0	0	F	0	H	0	0	0	0	M	0	P	0	R	0	0	0	0	0	0	0	Z	0	b	0	
F23	10	A	A	C	D	0	F	0	H	I	0	K	0	F	0	P	0	R	D	0	K	0	W	0	0	H	0	R	0	
F24	10	A	B	A	0	E	F	0	H	0	J	0	L	M	0	P	0	H	0	T	0	J	0	E	0	M	0	F	0	
F25	10	A	B	B	0	0	F	G	F	0	0	0	0	M	N	P	N	M	0	0	0	0	0	0	0	Y	Z	a	Z	Y
F26	10	A	B	C	D	0	F	0	H	3*	0	K	0	6*	0	P	0	R	B/2	0	16*	0	W	0	0	22*	0	2W	0	
F27	10	A	B	C	A/2	0	F	0	H	1*	0	H/2	0	M	0	P	0	R	12*	0	17*	0	R/2	0	0	Z	0	b	0	
F28	16	A	B	C	D	0	F	0	H	I	0	K	0	M	0	P	0	R	S	0	U	0	W	0	0	Z	0	b	0	
F29	16	A	B	C	0	E	F	0	H	0	J	0	L	M	0	P	0	R	0	T	0	V	0	X	0	Z	0	b	0	
F30	16	A	B	C	0	0	F	G	H	0	0	0	0	M	N	P	Q	R	0	0	0	0	0	0	0	Y	Z	a	b	c
F31	16	A	A	C	D	E	F	G	H	I	J	K	L	F	-J	P	Q	R	D	-G	K	-Q	W	X	-E	H	-L	R	-X	
F32	16	A	A	C	D	E	F	G	H	I	J	K	L	F	J	P	Q	R	D	G	K	Q	W	X	E	H	L	R	X	
F33	16	A	B	A	D	E	F	G	H	I	J	K	L	M	N	P	-K	H	S	T	-N	J	-G	E	-S	M	-I	F	-D	
F34	16	A	B	A	D	E	F	G	H	I	J	K	L	M	N	P	K	H	S	T	N	J	G	E	S	M	I	F	D	
F35	16	A	B	B	D	-D	F	G	F	I	J	-J	-I	M	N	P	N	M	S	T	U	-U	-T	-S	Y	Z	a	Z	Y	
F36	16	A	B	B	D	D	F	G	F	I	J	J	I	M	N	P	N	M	S	T	U	U	T	S	Y	Z	a	Z	Y	
F37	16	A	B	C	D	E	F	G	H	3*	J	K	L	6*	8*	P	Q	R	B/2	14*	16*	19*	W	X	20*	22*	23*	2W	2X	
F38	16	A	B	C	A/2	0	F	G	H	1*	G	H/2	0	M	N	P	Q	R	12*	15*	17*	Q	R/2	0	Y	Z	a	b	c	
F39	16	A	B	C	D	E	F	G	H	3*	J	K	L	6*	9*	P	Q	R	B/2	9*	16*	Q	W	X	0	22*	0	2W	0	
F40	16	A	B	C	A/2	E	F	E/2	H	1*	J	H/2	L	M	10*	P	L/2	R	12*	T	17*	V	R/2	X	21*	Z	24*	b	X/2	
F41	28	A	B	C	D	E	F	G	H	I	J	K	L	M	N	P	Q	R	S	T	U	V	W	X	Y	Z	a	b	c	

\* Notes. 1 -A/4+3F/2; 2 A/2-3D/2+3F/2; 3 B/20-3D/5+3F/2; 4 -2E/5+G; 5 A-2D+F; 6 B/5-2D/5+F; 7 -3E/5+G; 8 2E-5G+4J; 9 -G+2J; 10 -E/4+3J/2; 11 A-D; 12 A/2-5F/2+5M/2; 13 -E+G; 14 6E-15G+10J; 15 -G+2N; 16 -2K+3P; 17 -H/4+3P/2; 18 -L+Q; 19 -2L+3Q; 20 12E-30G+20J; 21 E/2-5J/2+5T/2; 22 -4K+6P; 23 -4L+6Q; 24 -L/4+3V/2.

Results

Calculations were carried out only for contra-variant coefficients; only these enter into the standard crystallographic least-squares procedure. The present tabulation covers all crystallographic settings of special positions for tensors up to rank 6 and all

cubic, hexagonal, tetragonal and trigonal settings for tensors up to rank 8. The orientation of the symmetry elements in the lattice must be identified (e.g. by inspection of *International Tables for X-ray Crystallography*, 1952) before the tables can be used. There is a key attributed to each site symmetry given in Table 1. The symmetry relations of this site

may be looked up under the corresponding entry in Table 2.\* The order of the coefficients corresponds exactly to the order of assignment of independency, *i.e.* in the case of linear relationships the independent parameters always enter on the left-hand side of the dependent ones. It is worth noting that three of the non-centrosymmetric groups have null third-rank tensors but one free coefficient as some higher level (rank 5 for 422, rank 7 for 622, rank 9 for 432).

It should also be remembered that, in a least-squares refinement of a non-centrosymmetric structure, one parameter corresponding to a non-zero entry for the point group has to be kept fixed (Hazell & Willis, 1978).

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\* The tables containing the symmetry restrictions of the seventh- and eighth-rank tensors have been deposited with the British Library Lending Division as Supplementary Publication No. SUP 38913 (11 pp.). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

was done at the Max-Planck Institut für Festkörperforschung in Stuttgart and at the Kristallographisches Institut in Freiburg.

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## The Solution of the One-Dimensional Sign Problem for Fourier Transforms

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### Abstract

An iterative procedure for the determination of the signs of scattering amplitudes is considered. It is assumed that the scattering density is a one-dimensional antisymmetric function with a limited range of definition. The convergence of the method to a rigorous solution is proved. The stability of the procedure with respect to various experimental errors is shown in model examples. The proof can be generalized for a one-dimensional phase determination of a continuous intensity distribution.

### Introduction

When non-crystalline objects are investigated by diffraction methods, the intensity of coherent scattering  $I(s)$  can often be measured as a continuous

function of scattering vector  $s$  (for instance, intensity distribution along layer lines for one-dimensionally periodic structures, intensity of small-angle scattering). The restoration of the scattering density distribution frequently requires the solution of the phase problem. The latter is analogous to the phase problem in crystal-structure analysis and lies in finding the phases of scattering amplitudes  $A(s)$  when their moduli are known from the experimental intensities. In the present paper the case will be considered when the scattering density is a one-dimensional antisymmetric function, so that its connection with the scattering amplitude is given by the sine-Fourier transform

$$A(s) = \mathcal{F}_s[\rho(r)] = \int_0^{\infty} \rho(r) \sin sr \, dr \quad (1)$$

and  $A(s)$  is a real function.